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Coordination-Dependent Equation of State of Expanded Liquid Alkali Metals and Critical Point Properties

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LETTER

Coordination-Dependent Equation of State of Expanded Liquid Alkali Metals and Critical Point Properties

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To take account of the known structure of expanded heavy alkali metals, a coordination numberdependent equation of state is set up. It is then shown that the critical constants T_c and V_s are related by $T_c V_c^{1/3}$ = constant, while the compressibility ratio is no longer constant through the alkali series, as required by experiment.

Key words: Neutron scattering. plasma coupling parameter.

The neutron scattering experiments of Hensel and coworkers¹ on expanded liquid alkali metals Rb and Cs along the liquid-vapour coexistence curve have demonstrated that the liquid structure factors are characterized by (i) a remarkably constant near-neighbour distance and (ii) an almost linear decrease of coordination number c with decreasing liquid density.

The purpose of the present work is to propose a form for the equation of state which reflects the above variation of coordination number with the thermodynamic state. The proposal will then be confronted with experiments on the critical constants p_c , V_c and T_c of the alkali metal series ranging from Li to Cs.

Earlier theoretical work carried out independently by Chapman and March² and by Likalter³ has expressed the pressure p in the equation of state as the sum of two terms: namely an excluded volume contribution p_{ex} and a Coulomb term p_{ch} proportional to the plasma coupling parameter $\Gamma = (Z^2e^2n^{1/3})/k_BT$, where *Z* is the valence and *n* is the ionic density. The basic idea underlying the present treatment is to supplement these two terms by a third contribution to the pressure which is now coordination number-dependent. Thus we shall write for the pressure

$$
p = p_{ex} + p_{ch} + p_{coord}
$$
 (1)

where p_{coord} takes the form

$$
p_{\text{coord}} = f(c, T). \tag{2}
$$

The single further assumption will be made, before committing ourselves to quite specific models for the first two terms appearing in eqn *(2),* that

$$
c = c(p, T) \tag{3}
$$

which means, on combining Eqns. $(1)-(3)$, that we are no longer dealing with an explicit form for the pressure p , but rather an implicit equation for p , due to the dependence of p_{coord} on p .

We shall now follow Likalter³ in modelling the excluded volume and Coulomb terms by writing

$$
p = \frac{n(1+Z)k_B T}{1-nb} - \frac{1}{3} \alpha Z^2 e^2 n^{4/3} + f(c, T),
$$
\n(4)

where α is the Madelung constant for the ionic assembly. Forming the spinodal curve defined by

$$
\left(\frac{\partial p}{\partial n}\right)_T = 0\tag{5}
$$

we then obtain

$$
p = \frac{1}{9}(1 - 4nb)\alpha Z^2 e^2 n^{4/3} + f(c, T).
$$
 (6)

The important point to note here is that since p_{coord} does not depend explicitly on n , its pressure derivative does not enter the spinodal, since it will be multiplied by $(\partial p/\partial n)_T$ which is itself zero from Eqn. (5).

The critical point is located by imposing the condition $(\partial p/\partial n)_T=0$ on the spinodal curve. One immediate and, to us, remarkable consequence of the proposed equation of state (1) is then that it yields

$$
T_c V_c^{1/3} = \text{constant.} \tag{7}
$$

This result is required from the empirical analysis of Chapman and March' who gave the exponent in Eqn. (7) as 0.3. That the value $1/3$ followed from Likalter's work was pointed out very recently by Klein and March⁴. The specific results for the critical density and temperature are

$$
n_c = \frac{1}{7b} \tag{8}
$$

and

$$
T_c n_c^{-1/3} = \frac{16 \alpha Z^2 e^2}{49(1+Z)}.
$$
\n(9)

In addition to Eqn.(7), there is a large variation in the compressibility ratio *Z,* defined by

$$
Z_c = \frac{p_c V_c}{RT_c} \tag{10}
$$

through the alkali series, ranging from 0.06 for Li to 0.20 for Rb and Cs. That such a situation is possible within the form of the equation of state proposed here becomes clear by forming Z_c from Eqns. (6)-(9), the result being

$$
Z_c = \frac{7}{48}(1+Z)\left[1 + \frac{21}{\alpha Z^2 e^2} n_c^{-4/3} f(c_c, T_c)\right].
$$
 (11)

Evidently, without committing oneself to a form of the function f , the conclusion is simply that *2,* does not follow as **a** constant through the alkali series.

While the result **(7),** in agreement with experiment already lends strong support to the assumptions **(1)-(3),** in work to be published elsewhere with D. J. Klein on crystalline metals, it can be shown that Eqn. **(3)** follows from current glue models of interatomic forces^{5,6} at $T= 0$.

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